

# Interaction of galaxies with environment – magnetic draping

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in collaboration with

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# Outline

- 1 Magnetic draping
  - Mechanism
  - Observations
  - Physical insight
- 2 Spiral galaxies
  - Polarized radio ridges
  - Magnetic draping simulations
  - Draping and synchrotron emission
- 3 Conclusions

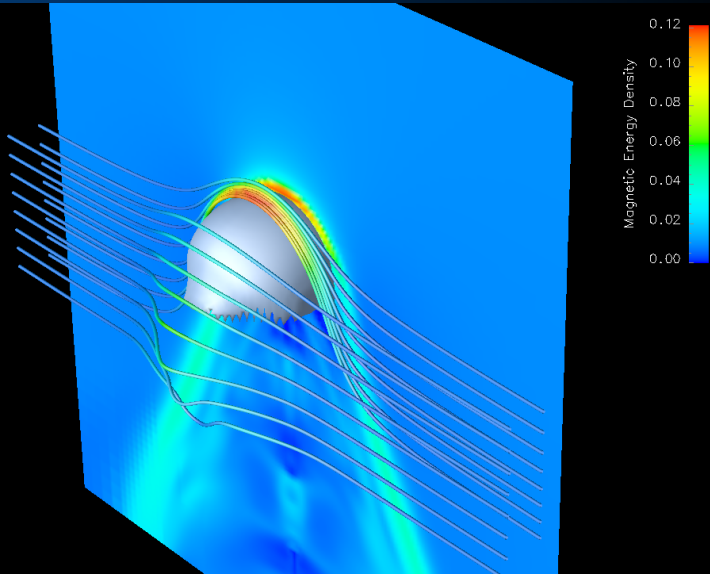


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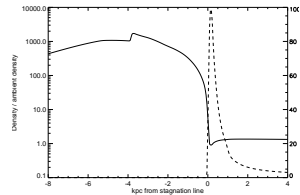


# What is magnetic draping?



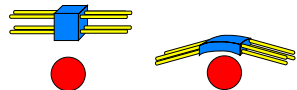
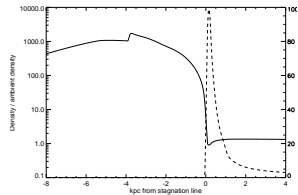
# What is magnetic draping?

- Is magnetic draping (MD) similar to ram pressure compression?
  - no density enhancement for MD
    - analytical solution of MD for incompressible flow
    - ideal MHD simulations (*right*)



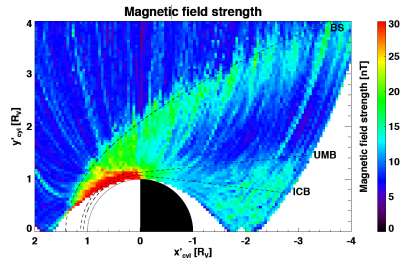
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→ no density enhancement for MD
  - analytical solution of MD for incompressible flow
  - ideal MHD simulations (*right*)
- Is magnetic flux still frozen into the plasma?  
yes, but plasma can also move along field lines while field lines get stuck at obstacle



# Draping of the interplanetary field over Venus

- Venus and Mars do not have a global magnetic field
- *Venus Express*: amplification of solar wind field by a factor  $\sim 6$  at the side facing the Sun



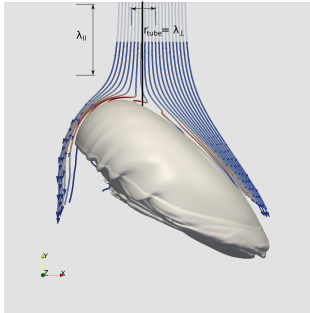
Guicking et al. (2010)

- draping of solar wind magnetic field around Venus/Mars leads to the **formation of magnetic pile-up region and the magneto tail**  
→ enhanced magnetic field strength in the planets' wake



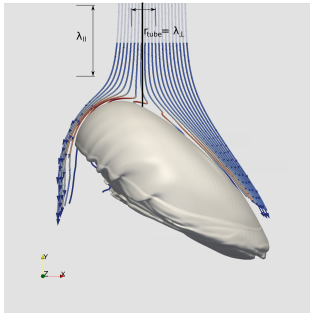


# Streamlines in the rest frame of the galaxy



- Stokes function  $p(s, \theta) = \sqrt{3sR} \sin \theta$   
→ critical impact parameter for  $\theta = \pi/2$ ,  $s = l_{\text{drap}}$ :  $p_{\text{cr}} = R/(2\mathcal{M}_A)$
- only those streamlines initially in a narrow tube of radius  $p_{\text{cr}} \simeq R/20 \simeq 1 \text{ kpc}$  from the stagnation line become part of the magnetic draping layer (color coded)  
→ constraints on  $\lambda_B$

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- the streamlines that do not intersect the tube get deflected away from the galaxy, become never part of the drape and eventually get accelerated (Bernoulli effect)
- note the kink feature in some draping-layer field lines due to back reaction as the solution changes from the hydrodynamic potential flow solution to that in the draped layer



# Conditions for magnetic draping

- **ambient plasma sufficiently ionized** such that flux freezing condition applies
- **super-Alfvénic motion** of a cloud through a weakly magnetized plasma:  $\mathcal{M}_A^2 = \beta\gamma\mathcal{M}^2/2 > 1$
- **magnetic coherence across the “cylinder of influence”:**

$$\frac{\lambda_B}{R} \gtrsim \frac{1}{\mathcal{M}_A} \sim 0.1 \times \left( \frac{\beta}{100} \right)^{-1/2} \text{ for sonic motions,}$$

Here  $R$  denotes the curvature radius of the working surface at the stagnation line.



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# Polarized synchrotron emission in a field spiral: M51

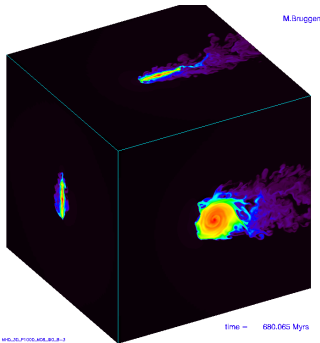


MPIfR Bonn and Hubble Heritage Team

- grand design 'whirlpool galaxy' (M51): optical star light superposed on radio contours
- polarized radio intensity follows the spiral pattern and is strongest in between the spiral arms
- the polarization 'B-vectors' are aligned with the spiral structure



# Ram-pressure stripping of cluster spirals

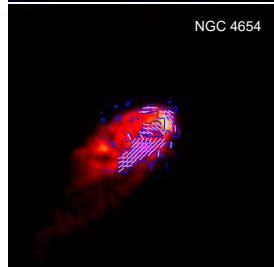
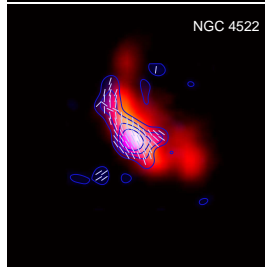
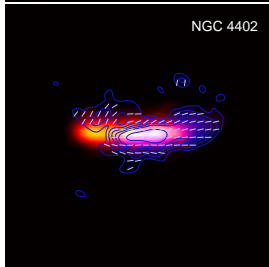
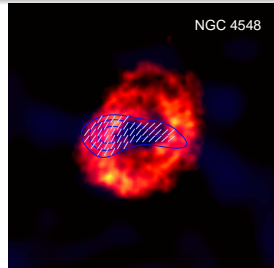
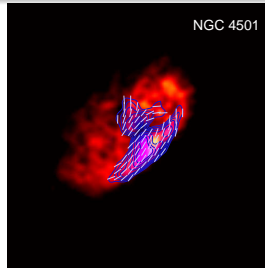
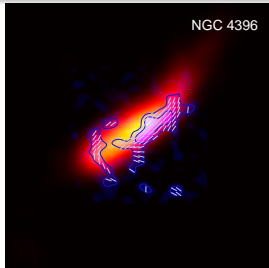


Brueggen (2008)

- 3D simulations show that the ram-pressure wind quickly strips the low-density gas in between spiral arms (Tonnesen & Bryan 2010)
- being flux-frozen into this dilute plasma, the large scale magnetic field will also be stripped

→ resulting radio emission should be unpolarized

# Polarized synchrotron ridges in Virgo spirals



Vollmer et al. (2007): 6 cm PI (contours) + B-vectors; Chung et al. (2009): HI (red)



# Observational evidence and model challenges

- asymmetric distributions of polarized intensity at the leading edge with extraplanar emission, sometimes also at the side
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$$n_{\text{compr}} \simeq n_{\text{icm}} v_{\text{gal}}^2 / c_{\text{ism}}^2 \simeq 1 \text{ cm}^{-3} \simeq \langle n_{\text{ism}} \rangle$$



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→ previous models that use ram-pressure compressed galactic magnetic fields fail to explain most of these points!



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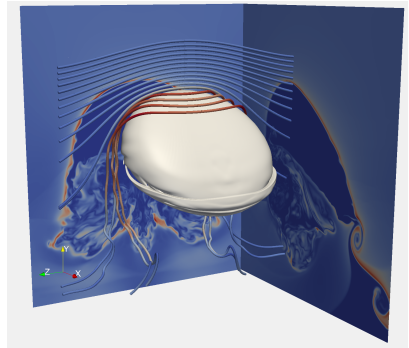
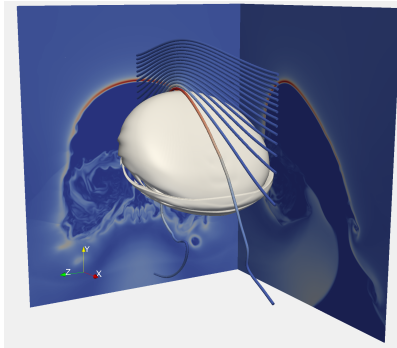
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→ need to consider the full MHD of the interaction spiral galaxy and magnetized ICM !



# Magnetic draping around a spiral galaxy

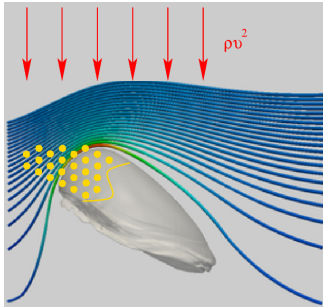


*Athena* simulations of spiral galaxies interacting with a uniform cluster magnetic field. There is a **sheath of strong field draped around the leading edge (shown in red)**.

C.P. & Dursi, 2010, Nature Phys.



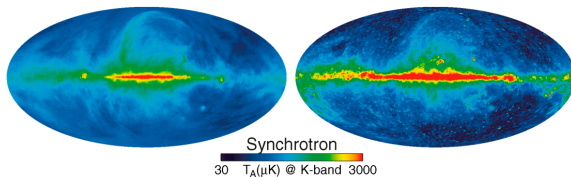
# Magnetic draping around a spiral galaxy – physics



- the galactic ISM is pushed back by the ram pressure wind  $\sim \rho v^2$
  - the stars are largely unaffected and lead the gas
  - the draping sheath is formed at the contact of galaxy/cluster wind
  - as stars become SN, their remnants accelerate CRes that populate the field lines in the draping layer
- 
- CRes are transported diffusively (along field lines) and advectively as field lines slip over the galaxy
  - CRes emit radio synchrotron radiation in the draped region, tracing out the field lines there → **coherent polarized emission at the galaxies' leading edges**



# Modeling the electron population



- typical SN rates imply a homogeneous CRe distribution (WMAP)
- FIR-radio correlation of Virgo spirals show comparable values to the solar circle → take CRe distribution of our Galaxy:

$$n_{\text{cre}} = C_0 e^{-(R-R_\odot)/h_R} e^{-|z|/h_z}$$

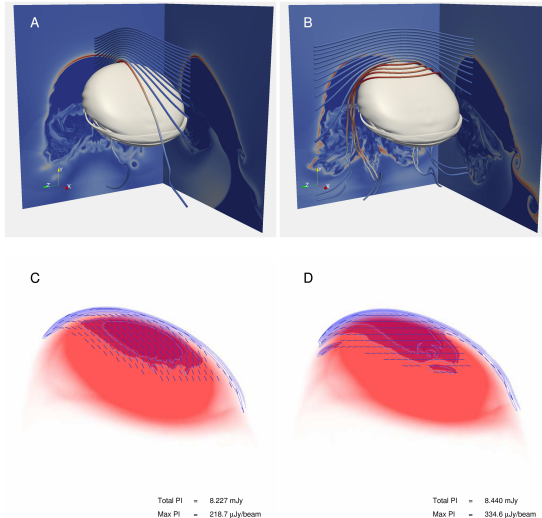
with normalization  $C_0 \simeq 10^{-4} \text{ cm}^{-3}$ ,  
scale heights  $h_R \simeq 8 \text{ kpc}$  and  $h_z \simeq 1 \text{ kpc}$  at Solar position

- truncate at contact of ISM-ICM, attach exp. CRe distribution  $\perp$  to contact surface with  $h_\perp \simeq 150 \text{ pc}$  (max. radius of Sedov phase)

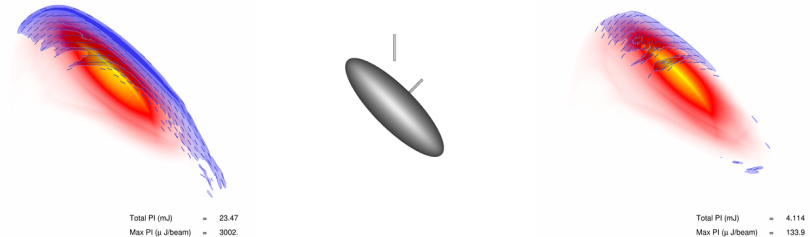


# Magnetic draping and polarized synchrotron emission

Synchrotron B-vectors reflect the upstream orientation of cluster magnetic fields



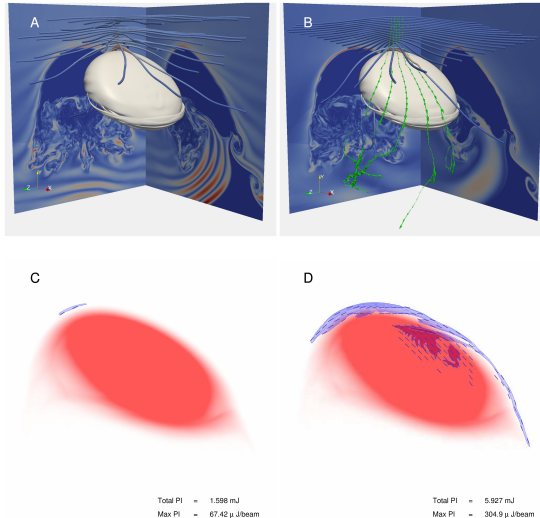
# Simulated polarized synchrotron emission



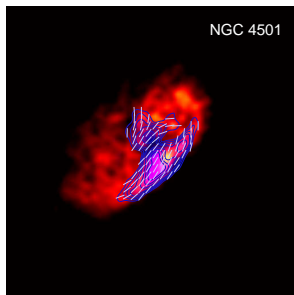
Movie of the simulated polarized synchrotron radiation viewed from various angles and with two field orientations.

# Magnetic draping of a helical B-field

(Non-)observation of polarization twist constrains magnetic coherence length



# Magnetic coherence scale estimate by radio ridges



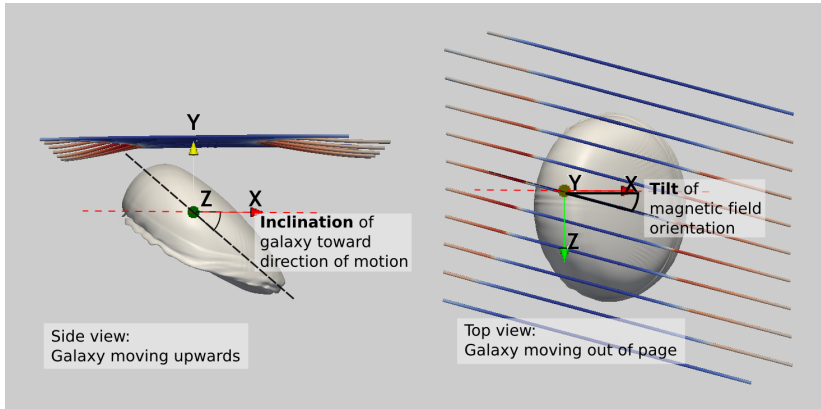
- observed polarised draping emission  
→ field coherence length  $\lambda_B$  is at least galaxy-sized
- if  $\lambda_B \sim 2R_{\text{gal}}$ , then the change of orientation of field vectors imprint as a change of the polarisation vectors along the vertical direction of the ridge showing a ‘polarisation-twist’
- the reduced speed of the boundary flow means that a small  $L_{\text{drape}}$  corresponds to a larger length scale of the unperturbed magnetic field ahead of the galaxy NGC 4501

$$L_{\text{coh}} \simeq \eta L_{\text{drape}} v_{\text{gal}} / v_{\text{drape}} = \eta \tau_{\text{syn}} v_{\text{gal}} > 100 \text{ kpc},$$

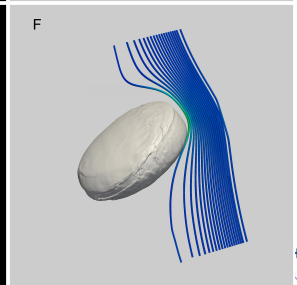
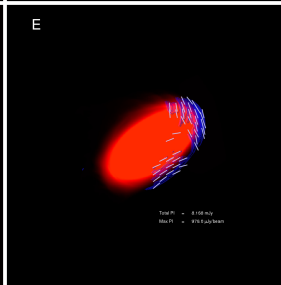
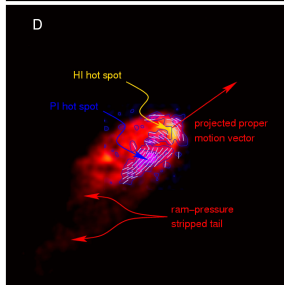
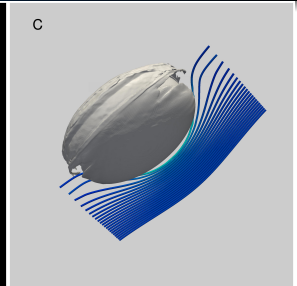
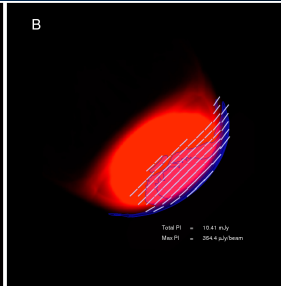
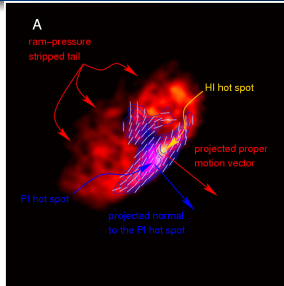
with  $\tau_{\text{syn}} \simeq 5 \times 10^7 \text{ yr}$ ,  $v_{\text{gal}} \simeq 1000 \text{ km/s}$ , and a geometric factor  $\eta \simeq 2$



# Varying galaxy inclination and magnetic tilt

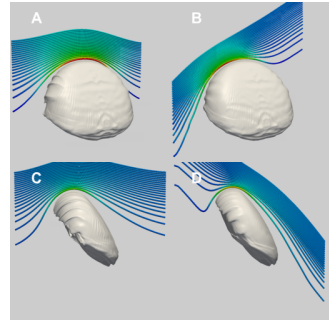


# Observations versus simulations

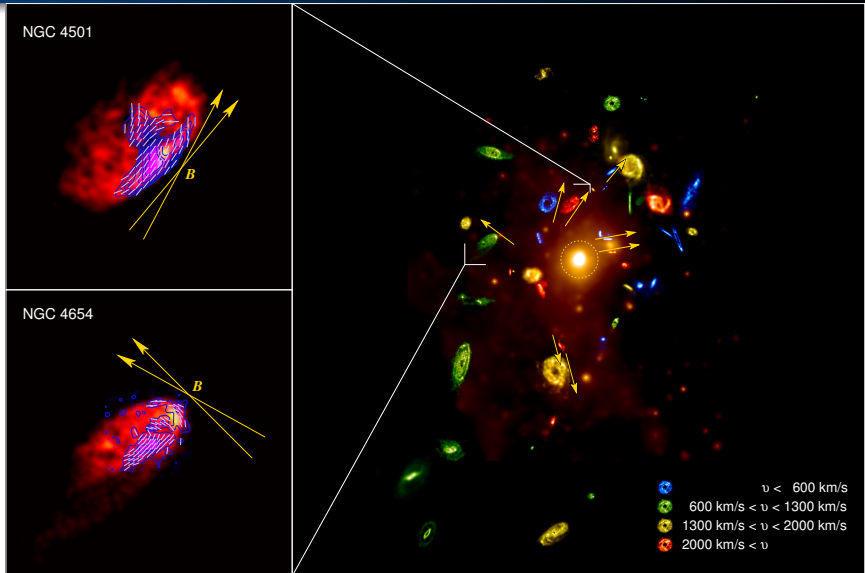


# Biases in inferring the field orientation

- uncertainties in estimating the 3D velocity:  $v_r$ , ram-pressure stripped gas visible in HI morphology  $\rightarrow \hat{\mathbf{v}}_t$
- *direction-of-motion asymmetry*: magnetic field components in the direction of motion bias the location of  $B_{\text{max, drape}}$  (figure to the right): draping is absent if  $\mathbf{B} \parallel \mathbf{v}_{\text{gal}}$
- *geometric bias*: polarized synchrotron emission only sensitive to traverse magnetic field  $B_t$  ( $\perp$  to LOS)  $\rightarrow$  maximum polarised intensity may bias the location of  $B_{\text{max, drape}}$  towards the location in the drape with large  $B_t$



# Mapping out the magnetic field in Virgo



# Discussion of radial field geometry

- The alignment of the field in the plane of the sky is **significantly more radial than expected from random chance**. Considering the sum of deviations from radial alignment gives a chance coincidence of less than 1.7% ( $\sim 2.2 \sigma$ ).
- For the **three nearby galaxy pairs** in the data set, **all have very similar field orientations**.

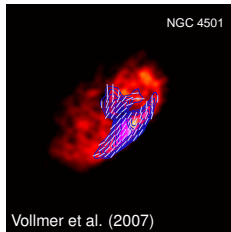
→ Which effect causes this field geometry?

**Magneto-thermal instability?** (Parrish+2007, C.P.+2010)

**Radial infall?** (Ruszkowski+2010)



# Magnetic draping with LOFAR



- NGC 4501: 5 GHz polarized intensity
- lower frequency  
→ longer electron cooling time  
→ **larger magnetic drape!**
- length scale of draping sheath:

$$\gamma = \left( \frac{2\pi\nu_{\text{syn}}m_e c}{3eB} \right)^{1/2} \simeq 10^4 \left( \frac{\nu_{\text{syn}}}{5 \text{ GHz}} \right)^{1/2} \left( \frac{B}{7 \mu\text{G}} \right)^{-1/2},$$

$$\tau_{\text{syn}} = \frac{6\pi m_e c}{\sigma_T B^2 \gamma} \simeq 50 \text{ Myr} \left( \frac{\nu_{\text{syn}}}{5 \text{ GHz}} \right)^{-1/2} \left( \frac{B}{7 \mu\text{G}} \right)^{-3/2},$$

$$L_{\text{drape}} = \eta v_{\text{drape}} \tau_{\text{syn}} \simeq 10 \text{ kpc} \left( \frac{\nu_{\text{syn}}}{5 \text{ GHz}} \right)^{-1/2} \simeq \mathbf{60 \text{ kpc}} \left( \frac{\nu_{\text{syn}}}{150 \text{ MHz}} \right)^{-1/2},$$

with velocity in draping layer  $v_{\text{drape}} \simeq 100 \text{ km s}^{-1}$  and a geometric factor  $\eta \simeq 2$ .

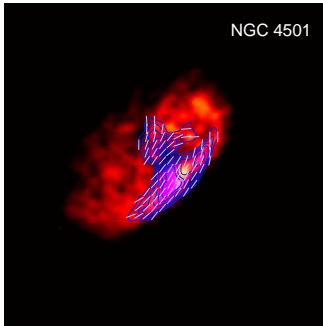


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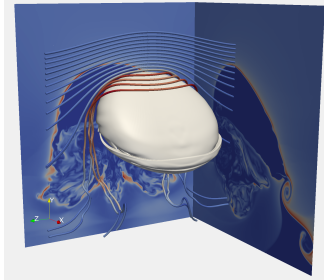


# Conclusions on magnetic draping around galaxies



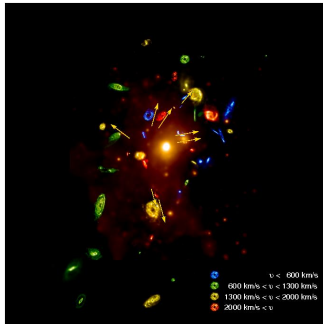
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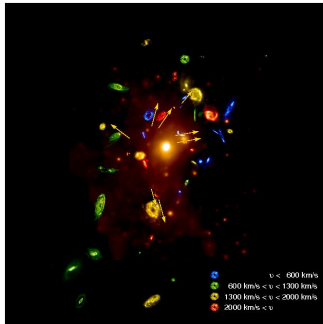
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- this represents a new tool for measuring the in situ 3D orientation and coherence scale of cluster magnetic fields
- application to the Virgo cluster shows that the magnetic field is preferentially aligned radially
- this finding implies efficient thermal conduction across clusters  
→ thermal cluster history & cluster cosmology
- great prospects for LOFAR observations:  
uncovering the history of the magnetic drape



# Literature for the talk

- Pfrommer & Dursi, 2010, Nature Phys., 6, 5206, *Detecting the orientation of magnetic fields in galaxy clusters*
- Dursi & Pfrommer, 2008, ApJ, 677, 993, *Draping of cluster magnetic fields over bullets and bubbles - morphology and dynamic effects*



# Magneto-thermal instability: the idea

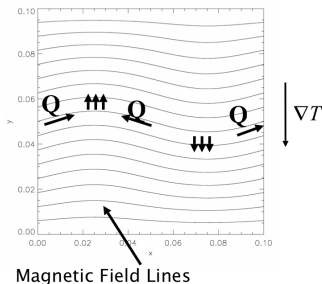


Figure from I. Parrish

Convective stability in a gravitational field:

- Classical Schwarzschild criterion:  
 $\frac{dS}{dz} > 0$
- long MFP, Balbus criterion:  $\frac{dT}{dz} > 0$
- **new instability causes field lines to reorient radially → efficient thermal conduction radially (close to Spitzer)**



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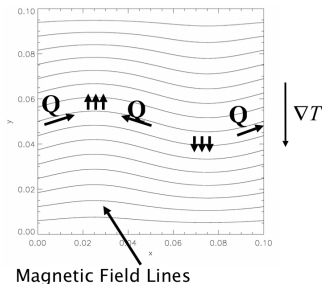


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The non-linear behavior of the MTI (Parrish & Stone 2007).

- **Adiabatic boundary conditions for  $T(r)$** : the instability can exhaust the source of free energy → isothermal profile
- **Fixed boundary conditions for  $T(r)$** : field lines stay preferentially radially aligned (35 deg mean deviation from radial)

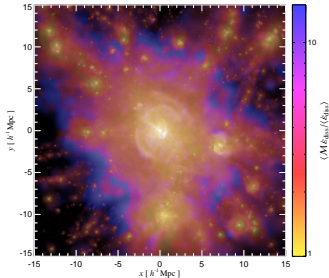


# Gravitational shock wave heating

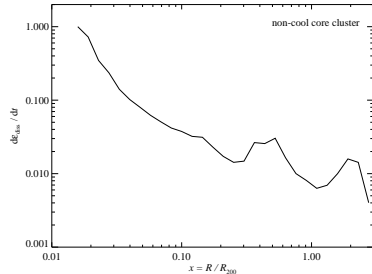
Observed temperature profile in clusters is decreasing outwards

→ heat also flows outwards along the radial magnetic field.

*How is the temperature profile maintained?* → gravitational heating



shock strengths weighted by dissipated energy



energy flux through shock surface

$$\dot{\epsilon}_{\text{diss}} / R^2 \sim \rho v^3$$

→ increase towards the center

